171. π -Orbital Interactions in *Möbius*-Type Molecules as Studied by Photoelectron Spectroscopy

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Summary

Some benzobarrelenes were studied by photoelectron (PE.) spectroscopy. The results can be interpreted by considering dominant through-space interaction of the π -molecular fragments in a *Möbius*-fashion. Some chemical features of the substrates are rationalized on the basis of these findings.

Introduction. – Barrelene (1), first synthesized by Zimmerman & Paufler in 1960 [1], clearly may be considered the archetypal Möbius molecule. Homoconjugative interaction among the three π -bonds must, by symmetry, involve an odd number of negative (antibonding) overlaps, regardless of how the π - and π^* -basis orbitals are defined in space. Detailed treatments of barrelene molecular orbitals have been presented elsewhere; the resulting orbital sequence [1] [2] is given in Figure 1. For a through-space interaction of magnitude β , a pair of degenerate antibonding MO's is lowered by β , with respect to the basis, while



Fig. 1. π -Molecular orbitals of barrelene derived from three interacting ethylene fragments

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the remaining two orbitals are raised by 2β . The overall splitting $(\Delta \varepsilon)$ is 3β in the bonding and antibonding manifold.

In a previous study [3] we have shown by PE. spectroscopy that this simple MO-treatment of barrelene is indeed quite valid, with two exceptions. First, there is apparently a significant contribution from hyperconjugative interaction with the lower-lying σ -skeleton for some π -levels, with the top one (the HOMO) excluded. This proved difficult to assess quantitatively, but is supported by recent MO-calculations by *Meyer & Pasternak* [4]. Second, there is a *Jahn-Teller* distortion in the first excited state of the radical cation 1⁺ which leads to a splitting of *ca*. 0.3 eV in the second ionization band.

In this communication, we continue our study of *Möbius* molecules within the usual framework of *Koopman*'s theorem and analyze the PE. spectra of a number of compounds in which one or more of the ethylenic π -orbitals have been replaced by a benzene ring [4]. These compounds include the well-known benzobarrelene (2), dibenzobarrelene (3), triptycene (4) series, in addition to appropriate model compounds (5-9).

We show here that orbital interactions in 2-4 can be understood by considering dominant through-space interaction of π -molecular fragments in a *Möbius*-fashion.



Basis molecular orbitals. – Our approach to the analysis of orbital interactions in benzobarrelenes first considers an appropriate collection of basis molecular orbitals [5]. We can then semi-quantitatively assess the relative through-space and through-bond contributions in model compounds 6, 7 and 8 and extend these to the fully homoconjugated systems 2, 3 and 4. Our initial collection of basis orbitals is simply the two highest-occupied benzene MO's, Ph_s and Ph_A and a bonding π -MO, π_s . The energies of these are taken as those of 5 [5] and 9[3] (Ph_s = -8.45 eV; Ph_A = -8.95 eV; π_s = -9.05 eV).



Neglecting overlap, through-space interactions of these orbitals can be calculated by first-order perturbation theory or, alternatively, by simple HMO-calculation. In the general case for homoconjugative interaction of two π -MO's, the resulting orbital splitting is given by *Equation 1*.

$$\Delta \varepsilon = 2\beta = 2(C_i C_j + C'_i C'_j)\beta^0.$$
⁽¹⁾

Here C_i/C_j and C'_i/C'_j are the orbital coefficients at the sites of nearest neighbor nonbonded interaction, β^0 is the resonance integral between the two basis atom (nonbonded) p-orbitals and $\Delta \varepsilon$ is the splitting of the two resulting π -orbitals. For the different possible types of basis-orbital interactions, we calculate the following $\Delta \varepsilon$ -values (Eqn. 2-5).

$$\pi_{\rm S}^1/\pi_{\rm S}^2 \qquad \Delta \varepsilon = 2\beta \tag{2}$$

$$Ph_{\rm S}^1/Ph_{\rm S}^2 \quad \Delta \varepsilon = \beta \tag{3}$$

$$Ph_{A}^{1}/Ph_{A}^{2} \ \Delta \varepsilon = 1/3\beta \tag{4}$$

$$Ph_S/\pi_S \quad \Delta \varepsilon = 2/\sqrt{2\beta}$$
 (5)

For the Ph_S/π_S -interaction, we here assume degeneracy; in practice, this interaction will decrease as the energy difference between Ph_S and π_S increases.

We next consider the symmetrical fully homoconjugated MO's in barrelene (1) and triptycene (4). Here, each basis π -orbital interacts twice with other orbitals. This leads us to the following results for the orbitals of interest, with interactions being uniquely determined by symmetry (Eqn. 6-8).

$$\pi_{\rm S}^1/\pi_{\rm S}^2/\pi_{\rm S}^3 \qquad \Delta \varepsilon = 3\beta \text{ for } \mathbf{1}$$
(6)

$$Ph_{S}^{1}/Ph_{S}^{2}/Ph_{S}^{3} \quad \Delta \varepsilon = 3/2\beta$$
for 4
(7)

$$Ph_{A}^{1}/Ph_{A}^{2}/Ph_{A}^{3} \ \Delta \varepsilon = 1/2\beta$$
(8)



Fig. 2. He(I)-PE. spectra of compounds 2, 3, 4, 6 and 7 (Due to different scales in the original spectra collected at different laboratories, they were redrawn using an Apple Graphics Tablet, the uniform scale being set by the spectrum of 4. As a result, the upper four spectra are smoothed with respect to noise and/or vibrational fine structure)

A critical feature in this type of analysis is the chosen value of the interaction parameter β ; in practice this usually is determined empirically. Furthermore, we need to distinguish between fully and partly homoconjugated systems. The former case is best represented by barrelene (1) itself with $\beta \approx -0.54$ eV, while β for the latter is taken as -0.29 eV. These values are easily derived from the experimentally observed orbital splittings in barrelene (1) (1.61 eV = -3β) [3] and dihydrobarrelene (8) (0.58 eV = -2β) [6]. The difference between these two values has been ascribed largely to differences in geometry [6].

Our analysis does not entirely ignore through-bond (hyperconjugative) effects since these are included to a certain degree in the empirically determined β -values. Furthermore, corrections for inductive effects based on appropriate model compounds will be made.

Results. – He(I)-PE. spectra were measured for the compounds 2, 3, 4, 6 and 7, preparations of which are described in [7]. The spectra are reproduced in *Figure 2* and the results summarized in the *Table*, together with the literature values for 1, 5, 8 and 9.

Compound		Vertical ionization potentials [eV]	
	1	$8.23, 9.65^{a}, 10.02^{a}$ [3]	
	2	8.12, 8.84, 9.45, 9.75, 11.5	
	3	8.0, 8.78 ^b), 9.50	
	4	7.9, 8.9 ^b)	
	5	8.45, 8.95 [5]	
	6	8.54, 8.94, 9.20	
	7	8.30, 8.80 ^b), 10.3	
	8	8.87, 9.45 [6]	
	9	9.05 [6]	

Table. Relevant vertical ionization potentials (eV) for barrelene derivatives measured by PE. spectroscopy

Discussion. – We consider first the PE. spectra of model compounds 6 and 7, which involve only one homoconjugative interaction. For 6, Ph_A should be essentially unchanged relative to model compound 5 (-8.95 eV), while through-space interaction of Ph_S and π_S will be small (*ca.* 0.1–0.15 eV) since these are strongly nondegenerate from a comparison of 5 and 9. This is essentially what is observed. We assign the three ionizations at 8.54, 8.94 and 9.20 eV as corresponding to the levels (Ph_S $-\pi_S$), Ph_A and (π_S +Ph_S), respectively. In this case, the lowering of π_S is primarily of inductive origin. Ph_S will be raised through-space and lowered inductively, with little net change.

For 7, we anticipate interaction between degenerate pairs of phenyl orbitals. From *Equation 3* and assuming $\beta = -0.29$ eV, we predict the $(Ph_s^1 - Ph_s^2)$ level to correspond to the ionization at 8.30 eV, in remarkably good agreement with experiment. The remaining three ionizations are predicted to occur at 8.60 $(Ph_s^1 + Ph_s^2)$, 8.9 $(Ph_A^1 - Ph_A^2)$ and 9.00 eV $(Ph_A^1 + Ph_A^2)$. In the observed spectrum, these are unresolved as a broad overlapping band. We now proceed to the analysis for the fully homoconjugated mono-, di- and tribenzobarrelenes 2, 3 and 4. For 2 we note two possible modes of analysis. First (*Model A*) we consider the interaction of basis orbitals of benzene (Ph_A and Ph_S) with those of 8: $(\pi_s^1 + \pi_s^1)$ and $(\pi_s^1 - \pi_s^2)$. This will emphasize interaction between the originally degenerate π -orbitals. The second model (*Model B*) considers three unmixed basis orbitals with preservation of the *Möbius* character of the orbital array.

The first possibility is shown in the interaction diagram (Fig. 3). Here $(\pi_s^1 + \pi_s^2)$ and Ph_A remain essentially unchanged due to their unique symmetries. Ph_s and $(\pi_s^1 - \pi_s^2)$ are of appropriate symmetry for interaction. From simple perturbation theory, we predict $\beta = -0.54$ eV, assuming degeneracy. Experimentally, Ph_s is raised *ca.* 0.41 eV and $(\pi_s^1 - \pi_s^2)$ lowered *ca.* 0.38 eV, in good agreement with this β -value whose absolute value is an upper bound for these splits due to actual non-degeneracy of the interacting orbitals.

Fig. 3. Orbital-interaction diagram for benzobarrelene

The alternative approach is to consider the three basis orbitals Ph_s, π_s^1 and π_s^2 prior to any interaction and mix these in a fashion which preserves the *Möbius* character of the array, according to the *secular determinant* 9 (in eV). Off-diagonal

	π_{S}^{1}	$\pi_{\mathbf{S}}^2$	Phs	
$\pi^1_{\mathbf{S}}$	$-9.25-\varepsilon$	+ 0.54	-0.38	(9)
$\pi_{s}^{\overline{2}}$	+0.54	$-9.25-\varepsilon$	-0.38	
Phs	-0.38	-0.38	$-8.53-\varepsilon$	

elements here are $\beta = -0.54$ eV from 1 (*i.e.* π_s^1/π_s^2) and $\beta = -0.38$ from Equation 5 (*i.e.* Ph_s/π_s). Diagonalization yields the solutions $\varepsilon = -9.79$, -9.14 and -8.02 eV. Model A seems to better reproduce experimental results, while Model B is

heuristically more satisfying, since it considers the full molecular symmetry.

In like manner, *Möbius* interaction in 3 can be described as in the secular determinant 10. Diagonalization yields the solutions $\varepsilon = -9.48$, -8.72 and

	Ph ¹ _S	Ph _S ²	π_{s}	
Ph_S^1	$-8.45-\varepsilon$	+ 0.27	-0.38	(10)
Ph_{S}^{2}	+0.27	$-8.45-\varepsilon$	-0.38	
π_{S}	-0.38	-0.38	$-9.25-\varepsilon$	

-7.95 eV. The observed spectrum displays well-resolved bands at 8.00 and 9.50 eV, almost exactly as predicted, and a broad unresolved system (maximum *ca.* 8.8 eV) which corresponds to the (Ph_S¹+Ph_S²)-level (which cannot by symmetry interact with π_S) and the two Ph_A-combinations (predicted at -8.90 and -9.00 eV).

Similar results are obtained through consideration of the alternative model; *i.e.* interaction of molecular orbitals of 7 and 9.

Finally, we consider triptycene (4) which possesses six separate *Möbius* arrays derived from symmetry combinations of the six degenerate sets of bonding and antibonding orbitals of the three benzene rings (*Fig. 4*). For the Ph_S-orbitals, assuming *Equation 7* to be valid, we predict the $(Ph_S^1 - Ph_S^2 - Ph_S^3)$ -combination to be at -7.91 eV (*i.e.* -8.45 + 0.54 eV), in excellent agreement with the experimental value (-7.90 eV). The remaining two Ph_S- and the Ph_A-combinations (with presumed *Jahn-Teller* splittings) all are observed as a broad unresolved band at *ca*. 8.9 eV, in agreement with *Figure 4*. Thus, our model correctly predicts the energy of the highest-occupied molecular orbital in triptycene with surprising accuracy. Our results thus demonstrate that through-space π/π -interaction is indeed quite significant in 4 with a magnitude essentially the same as in the parent compound 1. Unfortunately little additional information can be derived from the complex spectrum of 4.

Chemical consequences of orbital interactions. We consider here briefly some chemical properties of benzobarrelenes which might be influenced by interaction of π -component orbitals.

Fig. 4. Orbital interaction diagram for triptycene, neglecting Jahn-Teller effects

Concerning ground state properties, 1 exhibits an exceptionally large heat of hydrogenation for the first double bond which has recently been shown conclusively to result from through-space π/π -repulsions [8]. Our analysis indicates that 4 ought to behave qualitatively the same as 1, which agrees with the nonclassical strain deduced from its heat of formation [9].

Isolated double bonds (e.g. in 9) fail to undergo reduction in metal/ammonia solutions [10]. If H₂O is used as protonating agent for the intermediate anion, then also benzene fails to be reduced. Nevertheless, 2, 3 and 4 are easily reduced under these conditions, the former two being attacked at their double bonds [10] [11]. Clearly, this reactivity is related to homoconjugative interaction among the virtual orbitals of the π -fragments which does follow from the present findings in the bonding manifold. The resulting low-lying acceptor orbitals provide for stabilization in the intermediate anions.

Another well-studied area is photochemical reactions. Compounds 1 [12], 2 [13], 3 [14] and 4 [15] are all well-known to undergo photoreactions involving π/π -interaction and subsequent bond formation. While we do not expect a quantitative correlation of interaction with reactivity, it is nevertheless reasonable to suppose that these factors are related. Our results also help to explain the preference of benzobarrelene (2) for benzo-vinyl rather than vinyl-vinyl singlet excited state bridging. HOMO \rightarrow LUMO excitation in benzobarrelene should involve the transition given in Equation 13.

$$Ph_{S} - \pi_{S}^{1} - \pi_{S}^{2} \rightarrow Ph_{A}^{*} + \pi_{A}^{1*} - \pi_{A}^{2*}$$
 (13)

This transition promotes an electron from an orbital which is both benzo-vinyl and vinyl-vinyl antibonding to one which is benzo-vinyl bonding and vinyl-vinyl antibonding. Assuming this configuration to be dominant in the excited state wave function benzo-vinyl bonding should be the preferred mode of reaction, as was indeed confirmed by deuterium labeling [13]. We note in this context that such through-space orbital interactions were suggested in explaining the UV. spectra of heterocyclic analogs of benzobarrelenes [16].

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